

# The muon $g-2$ $\iff$ $\Delta\alpha$ connection

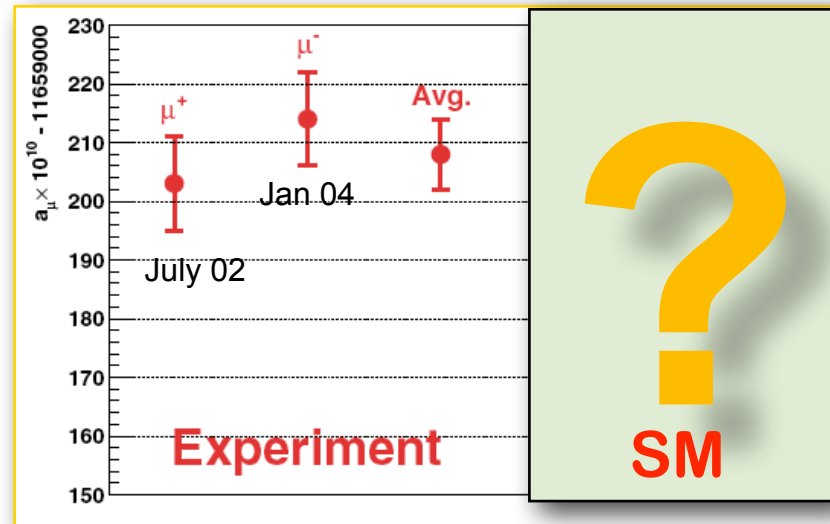
Massimo Passera  
INFN Padova

Joint Particle Physics Seminar  
University of Zurich and ETH Zurich  
09.11.2020

- **Muon g-2: recent theory progress**

- **Muon g-2  $\iff \Delta\alpha$  connection**

- **The MUonE project**



- **BNL 821:**  $a_\mu^{\text{EXP}} = (116592089 \pm 54_{\text{stat}} \pm 33_{\text{sys}}) \times 10^{-11}$  [0.5ppm].
- **Fermilab E989:** new muon g-2 experiment aims at  $\pm 16 \times 10^{-11}$  0.14ppm. First 3 data taking completed. Analysis of run 1 (~1xBNL) in progress. First result expected very soon with ~BNL precision.
- **J-PARC:** Muon g-2 proposal. Phase-1 with ~BNL precision.

# Muon $g-2$ : recent theory progress

White Paper of the Muon  $g-2$  Theory Initiative:  
arXiv:2006.04822

$$a_\mu^{\text{QED}} = (1/2)(\alpha/\pi) \quad \text{Schwinger 1948}$$

$$+ 0.765857426 (16) (\alpha/\pi)^2$$

Sommerfield; Petermann; Suura&Wichmann '57; Elend '66; MP '04

$$+ 24.05050988 (28) (\alpha/\pi)^3$$

Remiddi, Laporta, Barbieri ... ; Czarnecki, Skrzypek; MP '04;  
Friot, Greynat & de Rafael '05, Mohr, Taylor & Newell 2012

$$+ 130.8780 (60) (\alpha/\pi)^4$$

Kinoshita & Lindquist '81, ... , Kinoshita & Nio '04, '05;  
Aoyama, Hayakawa, Kinoshita & Nio, 2007, Kinoshita et al. 2012 & 2015;  
Steinhauser et al. 2013, 2015 & 2016 (all electron &  $\tau$  loops, analytic);  
Laporta, PLB 2017 (mass independent term). **COMPLETED!**

$$+ 750.86 (88) (\alpha/\pi)^5 \quad \text{COMPLETED!}$$

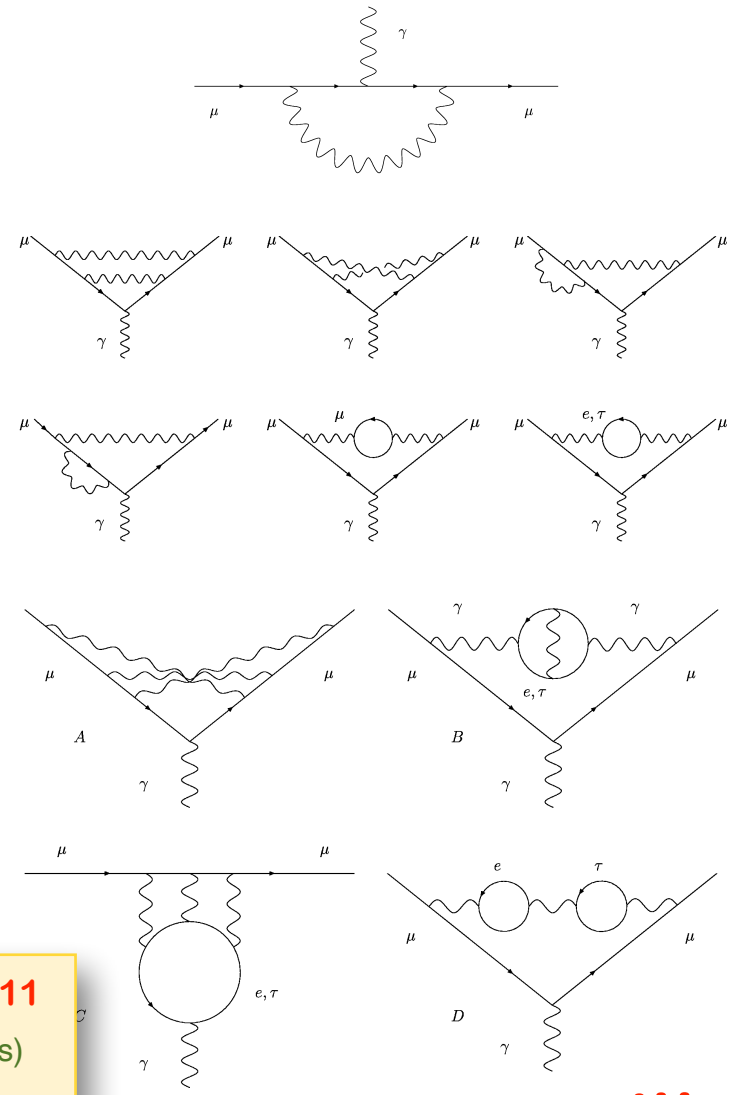
Kinoshita et al. '90, Yelkhovsky, Milstein, Starshenko, Laporta, ...  
Aoyama, Hayakawa, Kinoshita, Nio 2012, 2015, 2017 & 2019.  
Volkov 1909.08015:  $A_1^{(10)}$ [no lept loops] at variance, but negligible  $\Delta$ .

**Adding up, we get:**

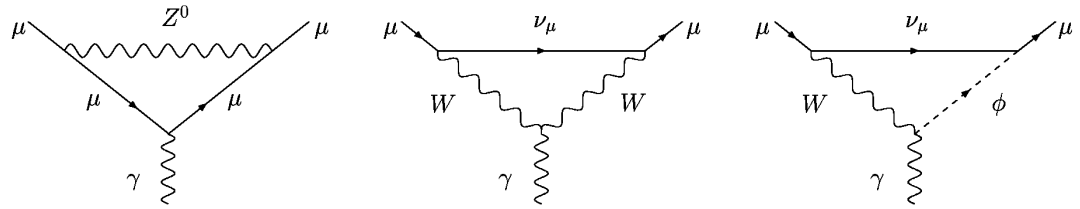
$$a_\mu^{\text{QED}} = 116584718.931 (19)(100)(23) \times 10^{-11}$$

from 4-loop & 5-loop coeffs unc. ↙ 6-loop ↘ from  $\alpha(\text{Cs})$

$$\alpha = 1/137.035999046(27) [0.2\text{ppb}] \quad 2018$$



● One-loop term:



$$a_{\mu}^{\text{EW}}(1\text{-loop}) = \frac{5G_{\mu}m_{\mu}^2}{24\sqrt{2}\pi^2} \left[ 1 + \frac{1}{5} (1 - 4\sin^2\theta_W)^2 + O\left(\frac{m_{\mu}^2}{M_{Z,W,H}^2}\right) \right] \approx 195 \times 10^{-11}$$

1972: Jackiv, Weinberg; Bars, Yoshimura; Altarelli, Cabibbo, Maiani; Bardeen, Gastmans, Lautrup; Fujikawa, Lee, Sanda; Studenikin et al. '80s

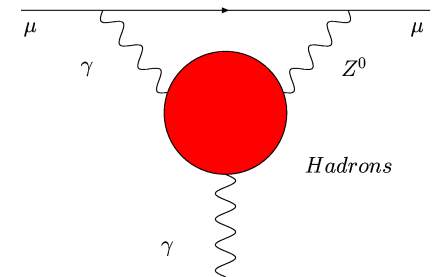
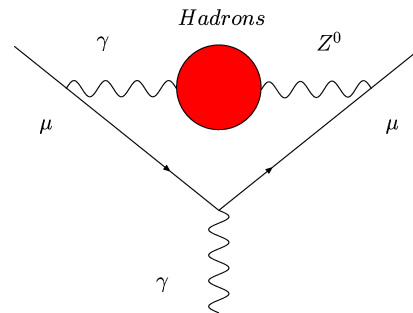
● One-loop plus higher-order terms:

$a_{\mu}^{\text{EW}} = 153.6 (1.0) \times 10^{-11}$

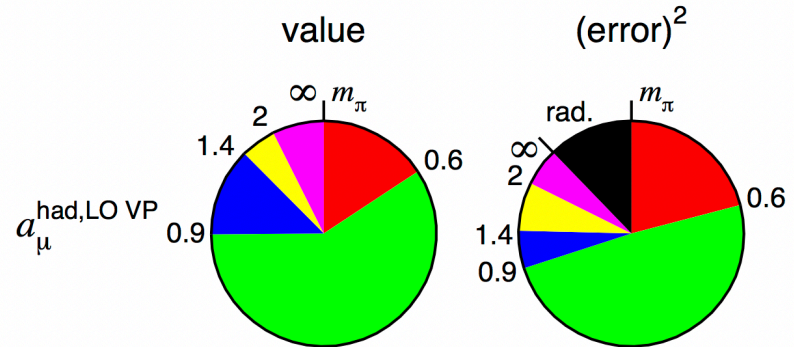
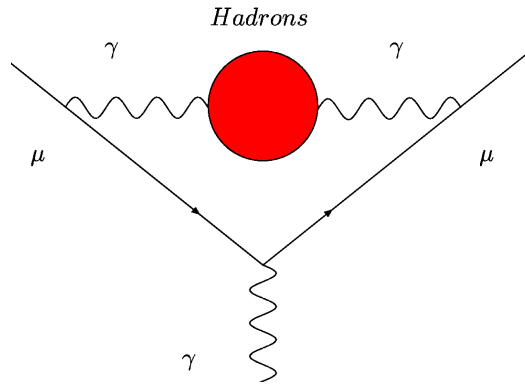
Kukhto et al. '92; Czarnecki, Krause, Marciano '95; Knecht, Peris, Perrottet, de Rafael '02; Czarnecki, Marciano and Vainshtein '02; Degrossi and Giudice '98; Heinemeyer, Stockinger, Weiglein '04; Gribouk and Czarnecki '05; Vainshtein '03; Gnendiger, Stockinger, Stockinger-Kim 2013, Ishikawa, Nakazawa, Yasui, 2019.

Hadronic loop uncertainties (and 3-loop nonleading logs).

Muon g-2 TI WP: arXiv:2006.04822



# The hadronic LO contribution



Keshavarzi, Nomura, Teubner 2018

$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)}$$

$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^\infty ds K(s) \sigma^{(0)}(s) = \frac{\alpha^2}{3\pi^2} \int_{4m_\pi^2}^\infty \frac{ds}{s} K(s) R(s)$$

$$a_\mu^{\text{HLO}} = 6895 (33) \times 10^{-11}$$

F. Jegerlehner, arXiv:1711.06089

$$= 6939 (40) \times 10^{-11}$$

Davier, Hoecker, Malaescu, Zhang, arXiv:1908.00921

$$= 6928 (24) \times 10^{-11}$$

Keshavarzi, Nomura, Teubner, arXiv:1911.00367

$$= 6931 (40) \times 10^{-11}$$

Muon g-2 TI WP: arXiv:2006.04822

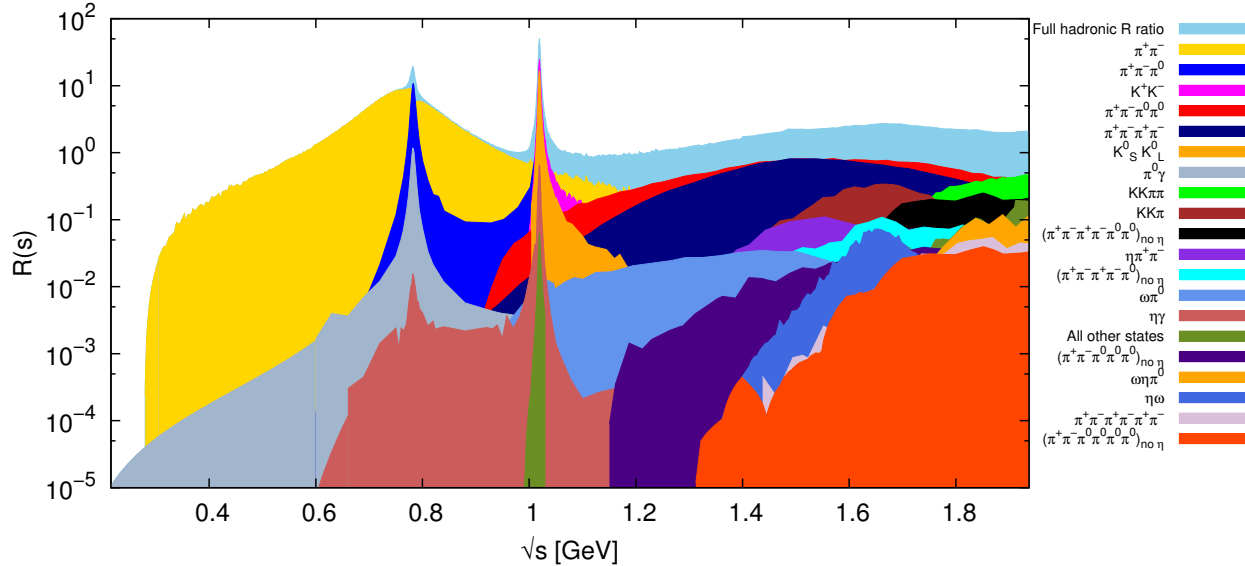


**Radiative Corrections to  $\sigma(s)$  are crucial.** S. Actis et al, Eur. Phys. J. C66 (2010) 585



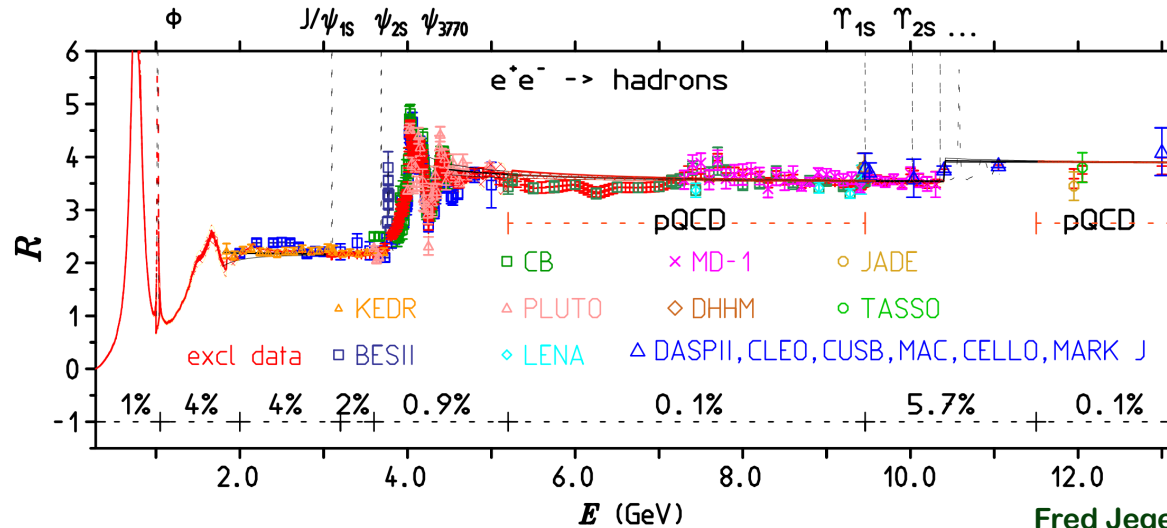
**Great progress in lattice QCD results. Recent BMW result with subpercent precision:**  
 $a_\mu^{\text{HLO}} = 7087(53) \times 10^{-11}$ . Tension with dispersive evaluations. S. Borsanyi et al. 2002.12347.

# The low-energy hadronic cross section



$$R(s) = \sigma(e^+e^- \rightarrow \text{hadrons}) / \frac{4\pi\alpha(s)^2}{3s}$$

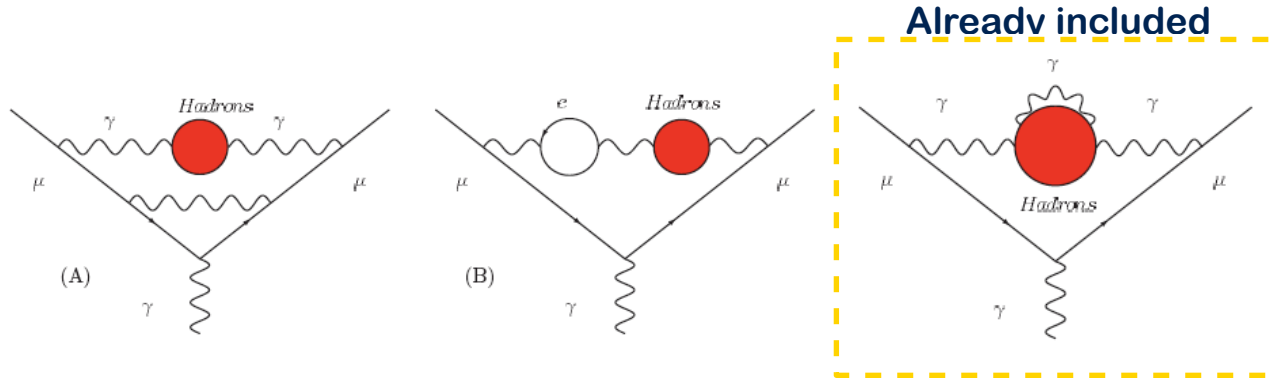
Keshavarzi, Nomura Teubner, PRD 2018



Fred Jegerlehner 1905.05078



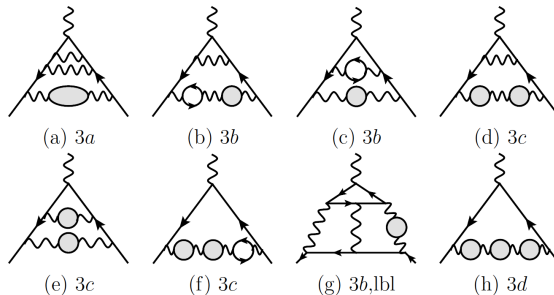
- $O(\alpha^3)$  contributions of diagrams containing HVP insertions:



$$a_{\mu}^{\text{HNLO}}(\nu p) = -98.3 (7) \times 10^{-11}$$

Krause '96; Keshavarzi, Nomura, Teubner 2019; Muon g-2 TI WP.

- $O(\alpha^4)$  contributions of diagrams containing HVP insertions:

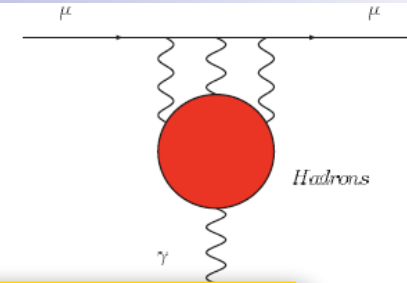


$$a_{\mu}^{\text{HNNLO}}(\nu p) = 12.4 (1) \times 10^{-11}$$

Kurz, Liu, Marquard, Steinhauser 2014

## ● HNLO light-by-light

This term had a troubled life! Nowadays:



$a_{\mu}^{\text{HNLO}}( b ) =$	$+ 80 (40) \times 10^{-11}$	Knecht & Nyffeler '02
	$= +136 (25) \times 10^{-11}$	Melnikov & Vainshtein '03
	$= +105 (26) \times 10^{-11}$	Prades, de Rafael, Vainshtein '09
	$= + 100 (29) \times 10^{-11}$	Jegerlehner, arXiv:1705.00263
	$= + 92 (19) \times 10^{-11}$	Muon g-2 TI WP, 2006.04822

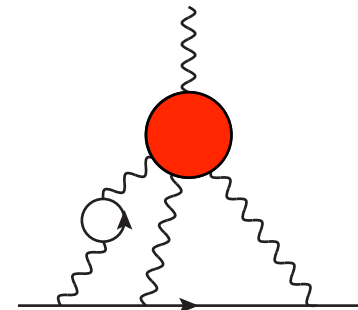
Significant improvements due to data-driven dispersive approach.

Great progress on the lattice. Recent RBC result:  $79(35) \times 10^{-11}$  arXiv:1911.08123

## ● HNNLO light-by-light

$$a_{\mu}^{\text{HNNLO}}(|b|) = 2 (1) \times 10^{-11}$$

Colangelo, Hoferichter, Nyffeler, MP, Stoffer 2014; Muon g-2 TI WP, 2006.04822



Comparing the SM prediction with the measured muon g-2 value:

$$a_{\mu}^{\text{EXP}} = 116592089 (63) \times 10^{-11}$$

BNL E821

$$a_{\mu}^{\text{SM}} = 116591810 (43) \times 10^{-11}$$

Muon g-2 TI

$$\Delta a_{\mu} = a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{SM}} = 279 (76) \times 10^{-11}$$

$3.7 \sigma$

# Muon $g-2 \iff \Delta\alpha$ connection

**Marciano, MP, Sirlin 2008 & 2010**  
**Keshavarzi, Marciano, MP, Sirlin 2020**

- Can  $\Delta a_\mu$  be due to **missing contributions** in the hadronic  $\sigma(s)$ ?
- An upward shift of  $\sigma(s)$  also induces an increase of  $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ .
- Consider:

$$a_\mu^{\text{HLO}} \rightarrow a = \int_{4m_\pi^2}^{s_u} ds f(s) \sigma(s), \quad f(s) = \frac{K(s)}{4\pi^3}, \quad s_u < M_Z^2,$$

$$\Delta\alpha_{\text{had}}^{(5)} \rightarrow b = \int_{4m_\pi^2}^{s_u} ds g(s) \sigma(s), \quad g(s) = \frac{M_Z^2}{(M_Z^2 - s)(4\alpha\pi^2)},$$

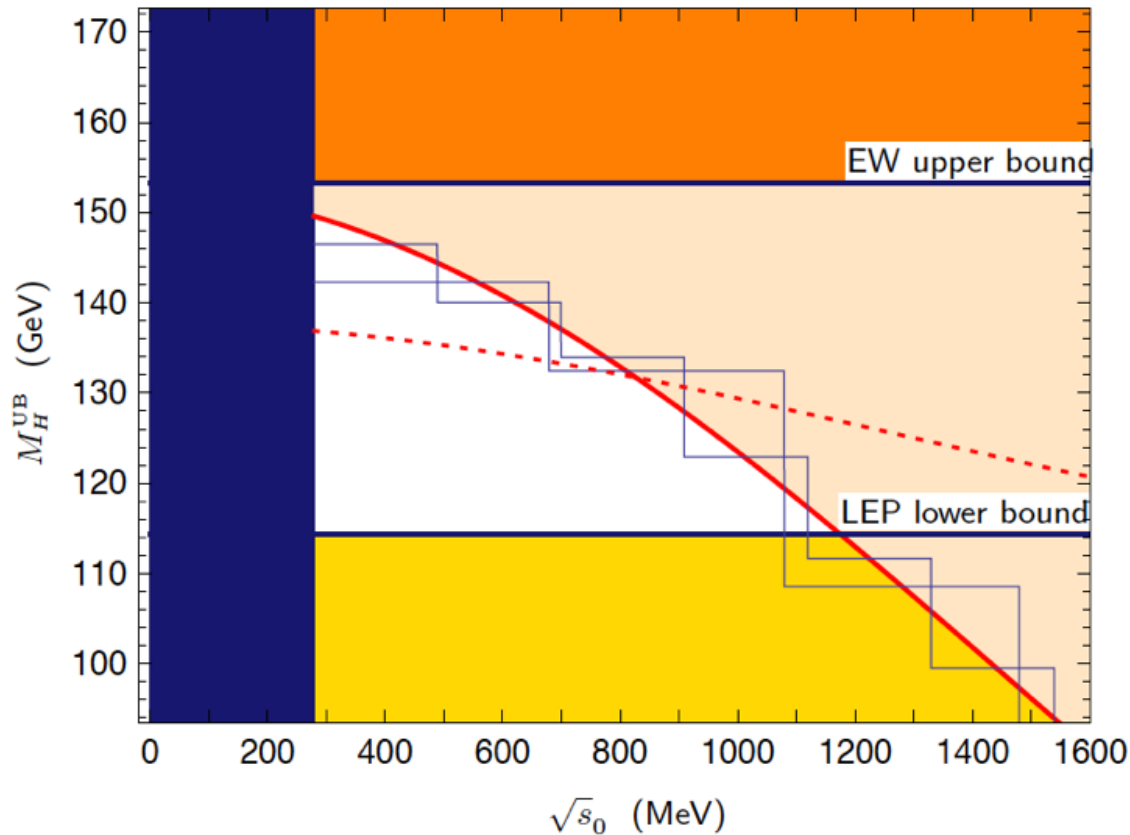
and the increase

$$\Delta\sigma(s) = \epsilon\sigma(s)$$

$\epsilon > 0$ , in the range:

$$\sqrt{s} \in [\sqrt{s_0} - \delta/2, \sqrt{s_0} + \delta/2] \quad \longrightarrow$$

How much does the  $M_H$  upper bound from the EW fit change when we shift up  $\sigma(s)$  by  $\Delta\sigma(s)$  [and thus  $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ ] to accommodate  $\Delta a_\mu$  ?

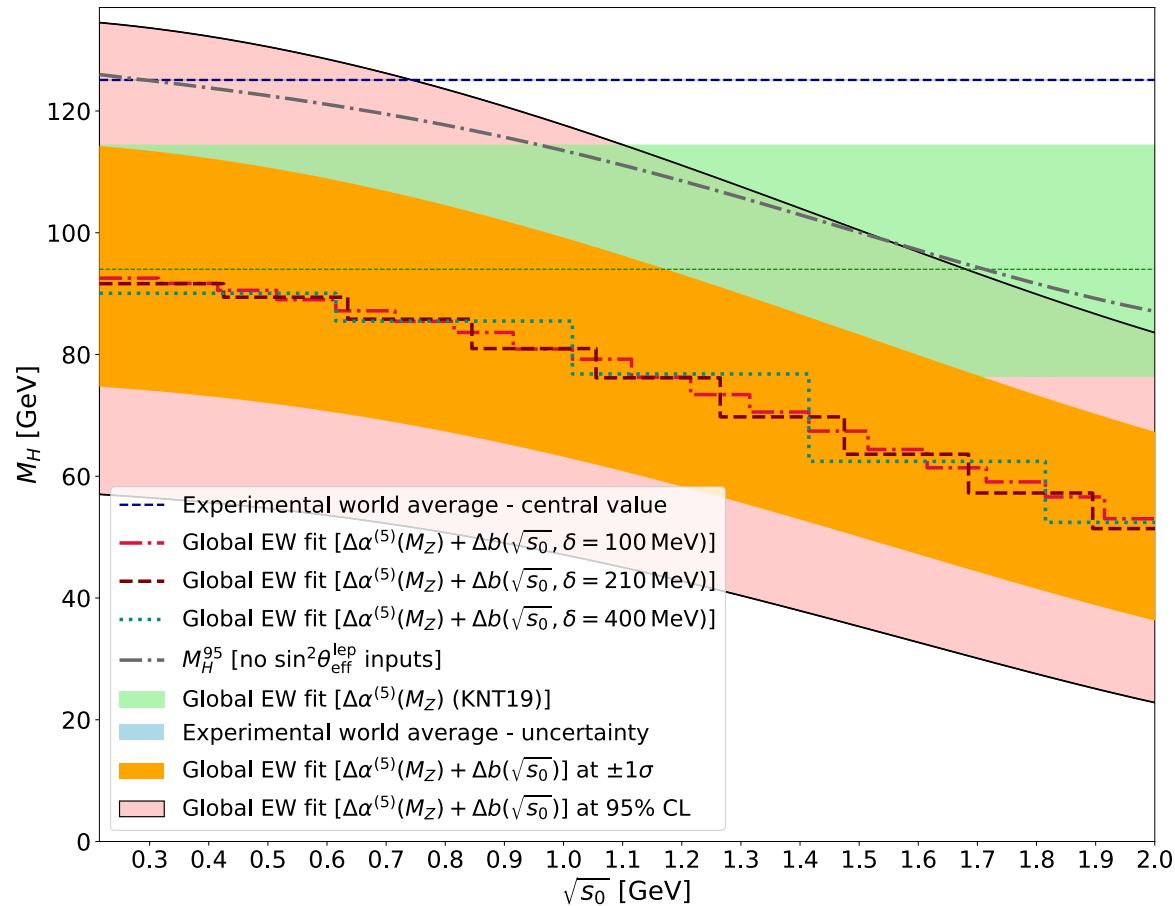


Marciano, MP, Sirlin, 2008 & 2010

Major update: Higgs discovered, improved EW observables ( $M_W$ ,  $\sin^2\theta$ ,  $M_{\text{top}}$ , ...), updates to  $\sigma(s)$ , theory improvements, global fit, ...

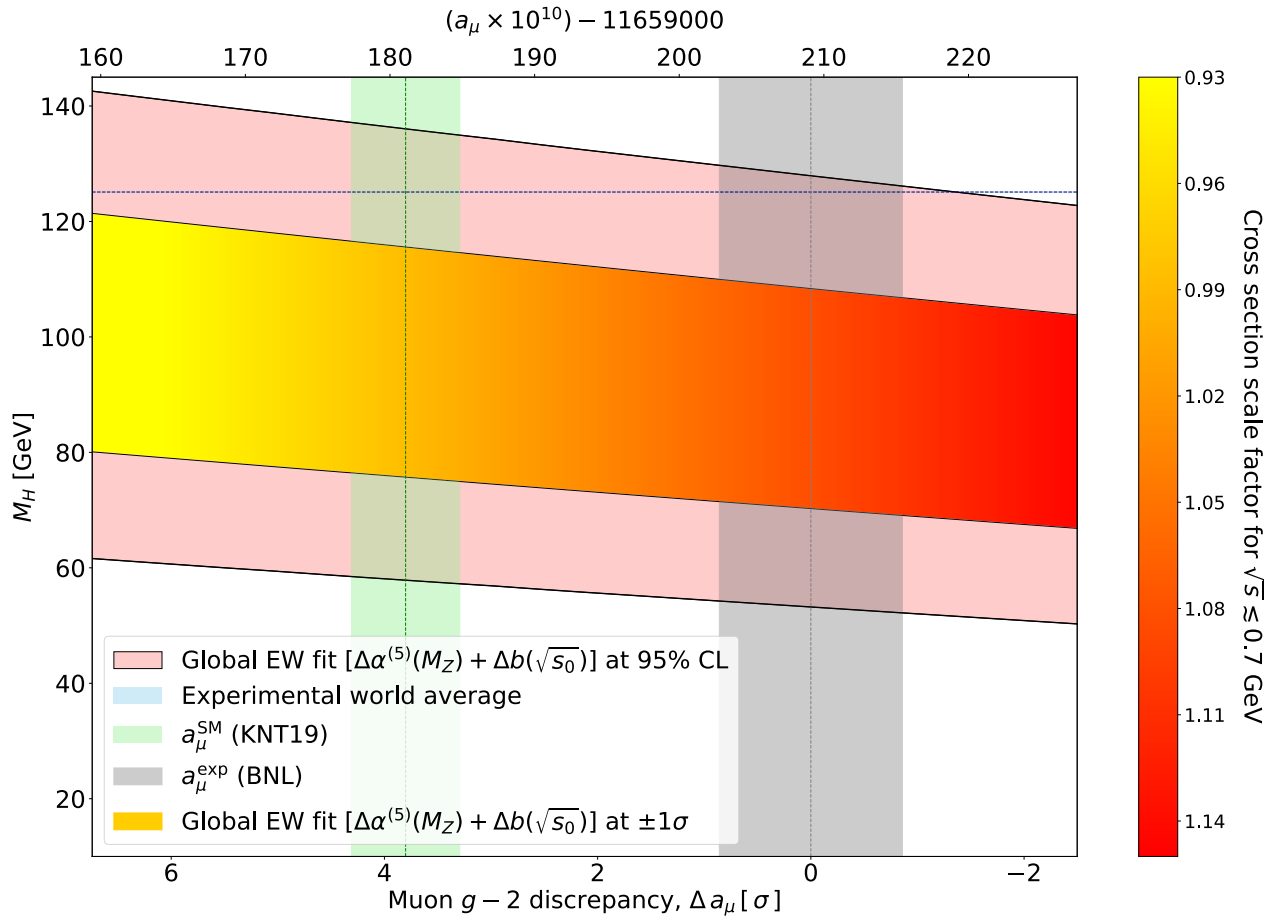
Parameter	Input value	Reference	Fit result	Result w/o input value
$M_W$ (GeV)	80.379(12)	[5]	80.359(3)	80.357(4)(5)
$M_H$ (GeV)	125.10(14)	[5]	125.10(14)	$94^{+20+6}_{-18-6}$
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) \times 10^4$	276.1(1.1)	[23]	275.8(1.1)	272.2(3.9)(1.2)
$m_t$ (GeV)	172.9(4)	[5]	173.0(4)	...
$\alpha_s(M_Z^2)$	0.1179(10)	[5]	0.1180(7)	...
$M_Z$ (GeV)	91.1876(21)	[5]	91.1883(20)	...
$\Gamma_Z$ (GeV)	2.4952(23)	[5]	2.4940(4)	...
$\Gamma_W$ (GeV)	2.085(42)	[5]	2.0903(4)	...
$\sigma_{\text{had}}^0$ (nb)	41.541(37)	[108]	41.490(4)	...
$R_l^0$	20.767(25)	[108]	20.732(4)	...
$R_c^0$	0.1721(30)	[108]	0.17222(8)	...
$R_b^0$	0.21629(66)	[108]	0.21581(8)	...
$\bar{m}_c$ (GeV)	1.27(2)	[5]	1.27(2)	...
$\bar{m}_b$ (GeV)	$4.18^{+0.03}_{-0.02}$	[5]	$4.18^{+0.03}_{-0.02}$	...
$A_{\text{FB}}^{0,l}$	0.0171(10)	[108]	0.01622(7)	...
$A_{\text{FB}}^{0,c}$	0.0707(35)	[108]	0.0737(2)	...
$A_{\text{FB}}^{0,b}$	0.0992(16)	[108]	0.1031(2)	...
$A_e$	0.1499(18)	[75,108]	0.1471(3)	...
$A_c$	0.670(27)	[108]	0.6679(2)	...
$A_b$	0.923(20)	[108]	0.93462(7)	...
$\sin^2\theta_{\text{eff}}^{\text{lep}}(Q_{\text{FB}})$	0.2324(12)	[108]	0.23152(4)	0.23152(4)(4)
$\sin^2\theta_{\text{eff}}^{\text{lep}}(\text{Had Coll})$	0.23140(23)	[100]	0.23152(4)	0.23152(4)(4)

Keshavarzi, Marciano, MP, Sirlin, PRD 2020 (using Gfitter)



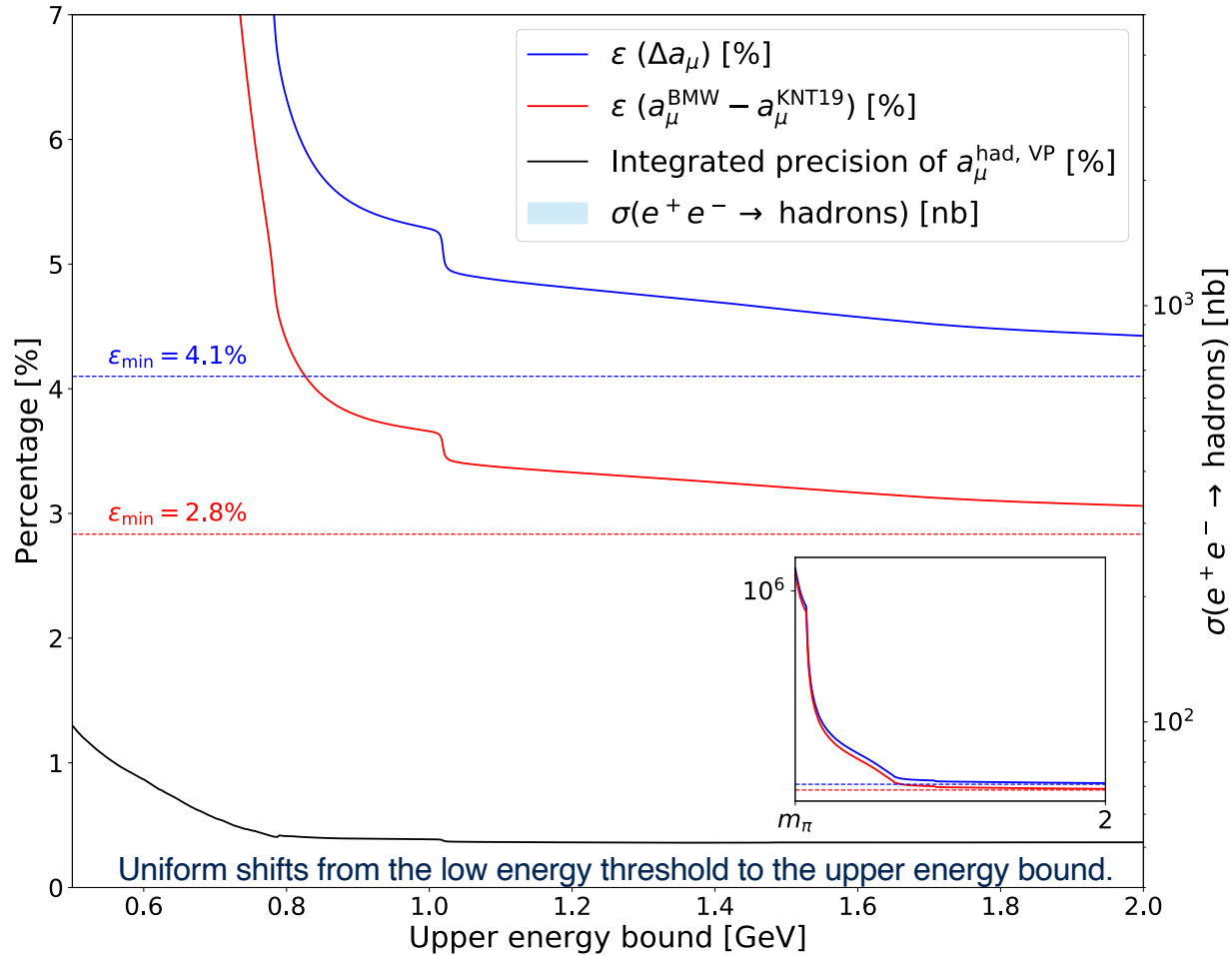
**Shifts  $\Delta\sigma(s)$  to fix  $\Delta a_\mu$  are possible,  
 but conflict with the EW fit if they occur above  $\sim 1$  GeV**





**Uniform scaling of  $\sigma(s)$  below  $\sim 0.7$  GeV? +9% required!**

Keshavarzi, Marciano, MP, Sirlin, PRD 2020



Shifts below 1 GeV conflict with the quoted exp. precision of  $\sigma(s)$

Keshavarzi, Marciano, MP, Sirlin, PRD 2020

# What happens to the electron $g-2$ ?

- The 2008 measurement of the electron g-2 is:

$$a_e^{\text{EXP}} = 11596521807.3 (2.8) \times 10^{-13} \quad \text{Hanneke et al, PRL100 (2008) 120801}$$

vs. old (factor of 15 improvement,  $1.8\sigma$  difference):

$$a_e^{\text{EXP}} = 11596521883 (42) \times 10^{-13} \quad \text{Van Dyck et al, PRL59 (1987) 26}$$

- Equate  $a_e^{\text{SM}}(\alpha) = a_e^{\text{EXP}}$  → “g<sub>e</sub>-2” determination of alpha:

$$\alpha^{-1} = 137.035\,999\,151 (33) \quad [0.24 \text{ ppb}]$$

- The best determination of  $\alpha$  is obtained via atomic interferometry:

$$\alpha^{-1} = 137.036\,999\,046 (27) [0.20 \text{ ppb}] \quad \text{Science 360 (2018) 191 (Cs)}$$

(was  $\alpha^{-1} = 137.035\,998\,995 (85) [0.62 \text{ ppb}]$  PRL106 (2011) & CODATA 2016 )

**2.5 sigma discrepancy**

- Using  $\alpha = 1/137.036\,999\,046\,(27)$  [Cs 2018], the SM prediction for the electron g-2 is:

$$a_e^{\text{SM}} = 115\,965\,218\,16.2\,(0.1)\,(0.1)\,(2.3) \times 10^{-13}$$

$\delta C_5^{\text{qed}}$      $\delta a_e^{\text{had}}$     from  $\delta\alpha$

- The (EXP – SM) difference is:

$$\Delta a_e = a_e^{\text{EXP}} - a_e^{\text{SM}} = -8.9\,(3.6) \times 10^{-13} [2.5\sigma]$$

NB: negative!

Note the negative sign. [QED 5-loop  $a_e^{\text{QED5}} = 4.6 \times 10^{-13}$ ]

- NP sensitivity limited only by the experimental errors in  $\alpha$  and  $a_e$ . May soon play a pivotal role in probing NP in the leptonic sector.

Giudice, Paradisi, MP 2012

- The present sensitivity is  $\delta\Delta a_e = 3.6 \times 10^{-13}$ , ie ( $10^{-13}$  units):

$$\underbrace{(0.1)_{\text{QED5}}, \quad (0.1)_{\text{HAD}}, \quad (2.3)_{\delta\alpha}, \quad (2.8)_{\delta a_e^{\text{EXP}}}}_{(0.1)_{\text{TH}}}$$

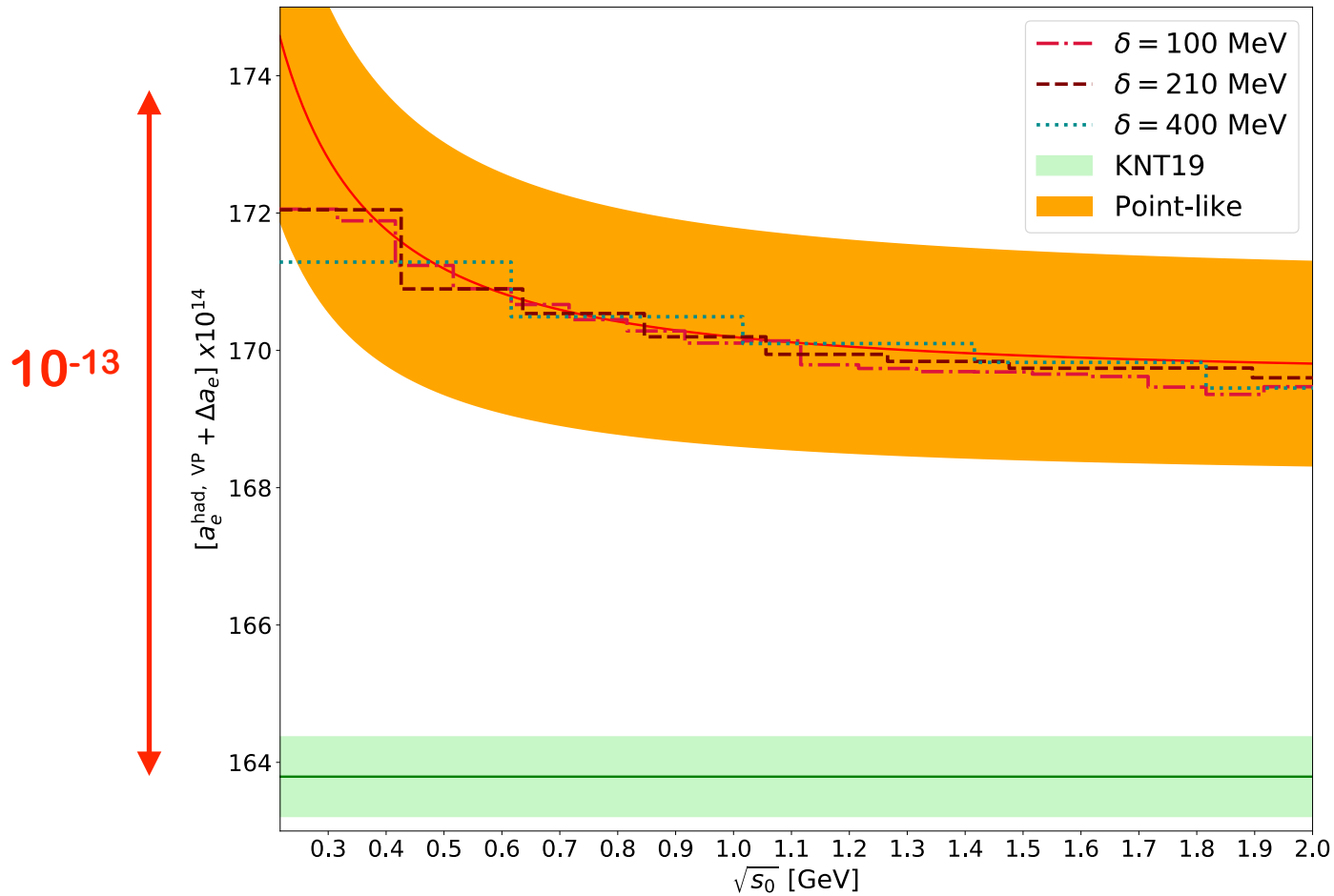
- The  $(g-2)_e$  exp. error may soon drop below  $10^{-13}$  and work is in progress to further reduce the error induced by  $\delta\alpha \rightarrow$

**sensitivity below  $10^{-13}$  may be reached with ongoing exp work**

- In a broad class of BSM theories, contributions to  $a_l$  scale as

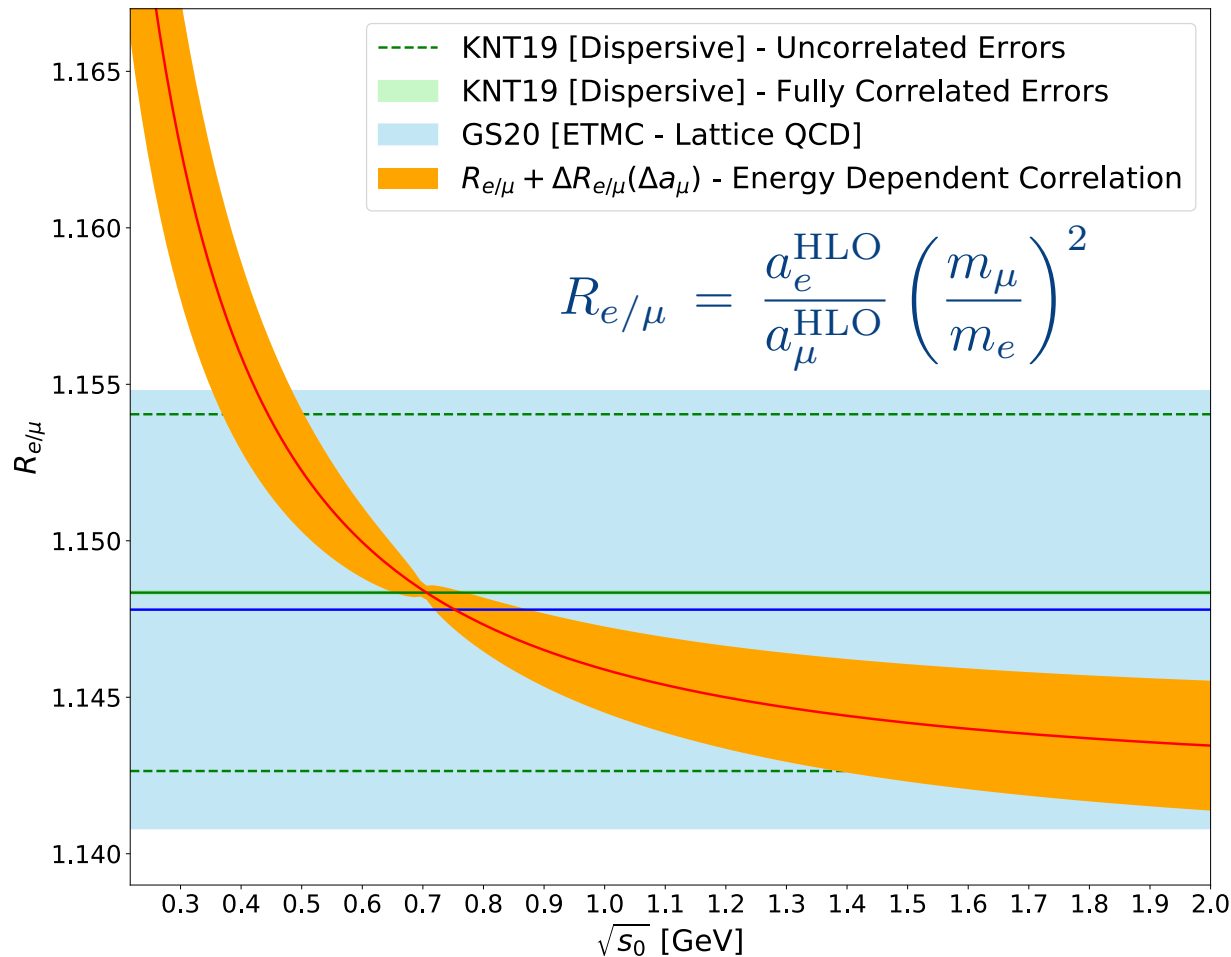
$$\frac{\Delta a_{\ell_i}}{\Delta a_{\ell_j}} = \left( \frac{m_{\ell_i}}{m_{\ell_j}} \right)^2 \quad \text{This Naive Scaling leads to:}$$

$$\Delta a_e = \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 0.7 \times 10^{-13}; \quad \Delta a_\tau = \left( \frac{\Delta a_\mu}{3 \times 10^{-9}} \right) 0.8 \times 10^{-6}$$



Shifts  $\Delta\sigma(s)$  to fix  $\Delta a_\mu$  slightly increase the  $|\Delta a_e| \sim 10^{-12}$  tension

Keshavarzi, Marciano, MP, Sirlin, PRD 2020



Good agreement between lattice [Giusti & Simula 2020] and KNT19.  
 Possible future bounds on very low energy shifts  $\Delta\sigma(s)$ ?

Keshavarzi, Marciano, MP, Sirlin, PRD 2020

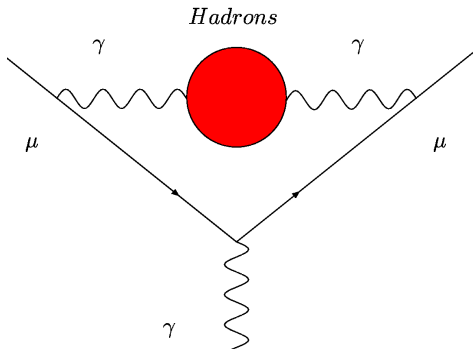


- Crivellin, Hoferichter, Manzari and Montull, “Hadronic vacuum polarization:  $(g-2)_\mu$  versus global electroweak fits,” PRL125 (2020) 091801 [arXiv:2003.04886].
- Eduardo de Rafael, “On Constraints Between  $\Delta\alpha_{\text{had}}(M_Z^2)$  and  $(g_\mu-2)_{\text{HVP}}$ ,” arXiv:2006.13880.
- Malaescu and Schott, “Impact of correlations between  $a_\mu$  and  $\alpha_{\text{QED}}$  on the EW fit”, arXiv:2008.08107.
- Colangelo, Hoferichter and Stoffer, “Constraints on the two-pion contribution to hadronic vacuum polarization,” arXiv:2010.07943.

# The MUonE project



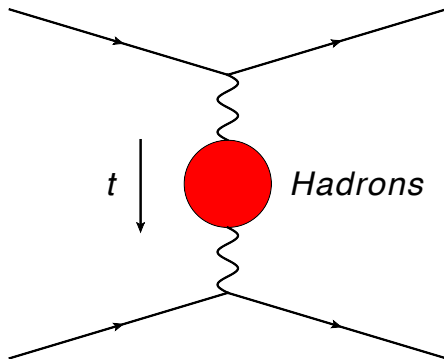
- The leading hadronic contribution  $a_\mu^{\text{HLO}}$  computed via the **timelike** formula:



$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} ds K(s) \sigma_{\text{had}}^0(s)$$

$$K(s) = \int_0^1 dx \frac{x^2 (1-x)}{x^2 + (1-x)(s/m_\mu^2)}$$

- Alternatively, simply exchanging the  $x$  and  $s$  integrations:



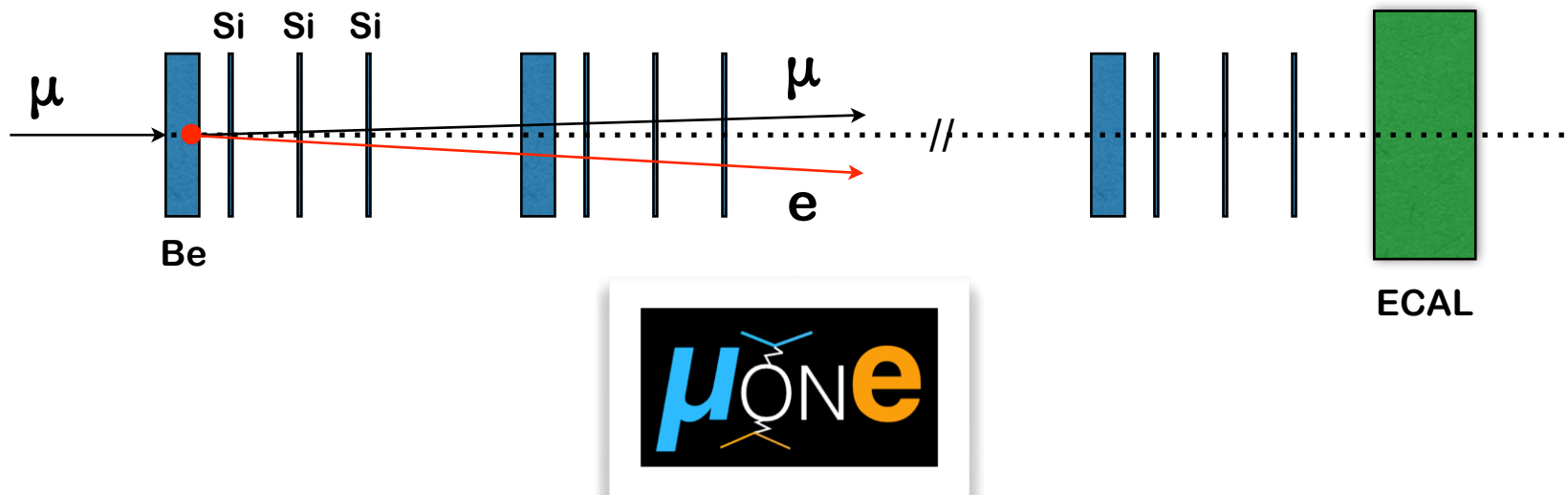
$$a_\mu^{\text{HLO}} = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{\text{had}}[t(x)]$$

$$t(x) = \frac{x^2 m_\mu^2}{x-1} < 0$$

Lautrup, Peterman, de Rafael, 1972

$\Delta\alpha_{\text{had}}(\mathbf{t})$  is the hadronic contribution to the running of  $\alpha$  in the **spacelike** region:  $a_\mu^{\text{HLO}}$  can be extracted from scattering data!

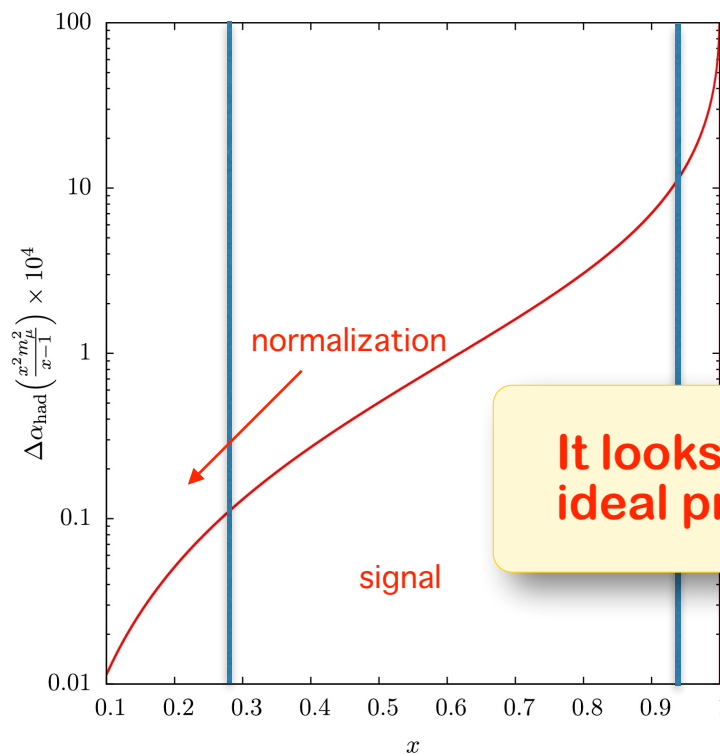
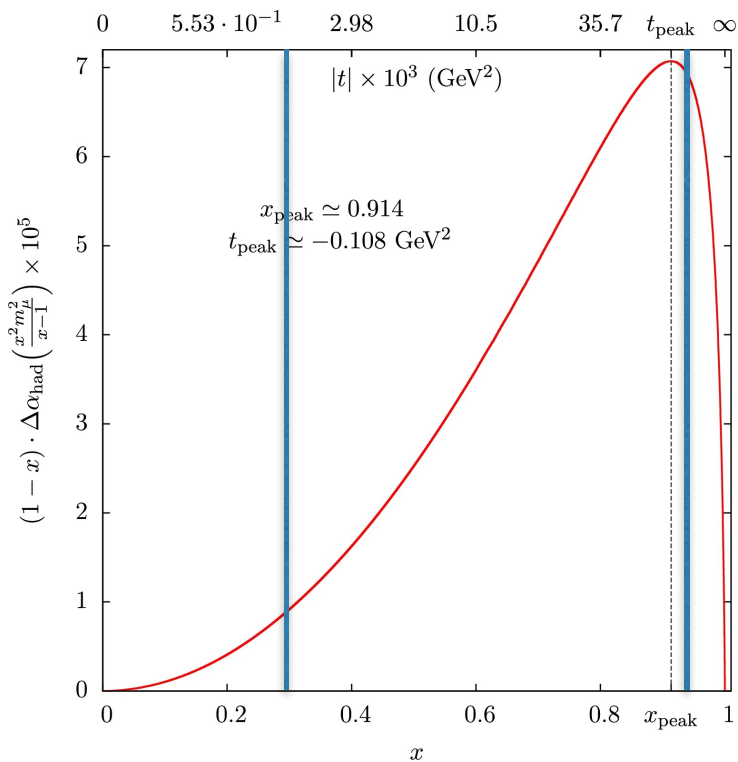
- $\Delta\alpha_{\text{had}}(t)$  can be measured via the **elastic scattering  $\mu e \rightarrow \mu e$** .
- We propose to scatter a 150 GeV muon beam, available at CERN's North Area, on a fixed electron target (Beryllium). Modular apparatus: each station has one layer of Beryllium (target) followed by several thin Silicon strip detectors.



Abbiendi, Carloni Calame, Marconi, Matteuzzi, Montagna,  
Nicosini, MP, Piccinini, Tenchini, Trentadue, Venanzoni

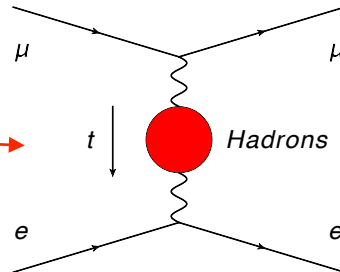
EPJC 2017 - arXiv:1609.08987

- For a 150 GeV muon beam, MUonE's scan region extends up to  $x=0.932$ , ie beyond the peak! (the peak is at  $x=0.914$ )
- The high-energy region inaccessible to MUonE contributes only 13% of  $a_\mu^{\text{HLO}}$  integral. It can be determined with timelike data and/or lattice QCD results. Recently obtained via lattice QCD. Giusti&Simula and Marinkovic'&Cardoso 2019



- **Statistics:** With CERN's 150 GeV muon beam M2 ( $1.3 \times 10^7 \mu/s$ ), incident on 40 15mm Be targets (total thickness 60cm), 2-3 years of data taking ( $2 \times 10^7$  s/yr)  $\rightarrow \mathcal{L}_{\text{int}} \sim 1.5 \times 10^7 \text{ nb}^{-1}$ .
- With this  $\mathcal{L}_{\text{int}}$  we estimate that measuring the shape of  $d\sigma/dt$  we can reach a statistical sensitivity of  **$\sim 0.3\%$  on  $a_{\mu}^{\text{HLO}}$** , ie  $\sim 20 \times 10^{-11}$ .
- **Systematic** effects must be known at  $\leq 10\text{ppm}$ !
- Test beams performed at CERN in 2017 & 2019 arXiv:1905.11677
- Lol submitted to CERN SPSC in 2019: Test run approved for 2021.
- Full-statistics run hopefully in 2022–24.

- To extract  $\Delta\alpha_{\text{had}}(t)$  from MUonE's measurement, the ratio of the SM cross sections in the signal and normalisation regions must be known at  $\lesssim 10\text{ppm}$ !



- **Fully differential fixed-order MC @ NLO ready** Pavia and PSI 2018-19
- **NNLO QED: MI for 2-loop box diagrams computed** Padova 2017-19
- **Two MC built including partial subsets of the NNLO QED corrections due to electron and muon radiation** Pavia and PSI 2020
- **NNLO hadronic effects computed** Padova and Siegen 2019
- **Extraction of the leading electron mass effects from the massless muon-electron scattering amplitudes** PSI 2019-2020
- **New physics at MUonE** Heidelberg&Washington and Padova 2020
- ...

Theory for muon-electron scattering @ 10 ppm:  
A report of the MUonE theory initiative. arXiv:2004.13663





2<sup>nd</sup> Workstop / Thinkstart: 4<sup>th</sup> – 7<sup>th</sup> Feb 2019

Theory for muon-electron scattering @ 10ppm

Y36 K08, Physik-Institut, University of Zurich

Organized by  
A. Signer & Y. Ulrich

**MUonE theory workshops: Padova 2017, Mainz 2018, Zurich 2019**

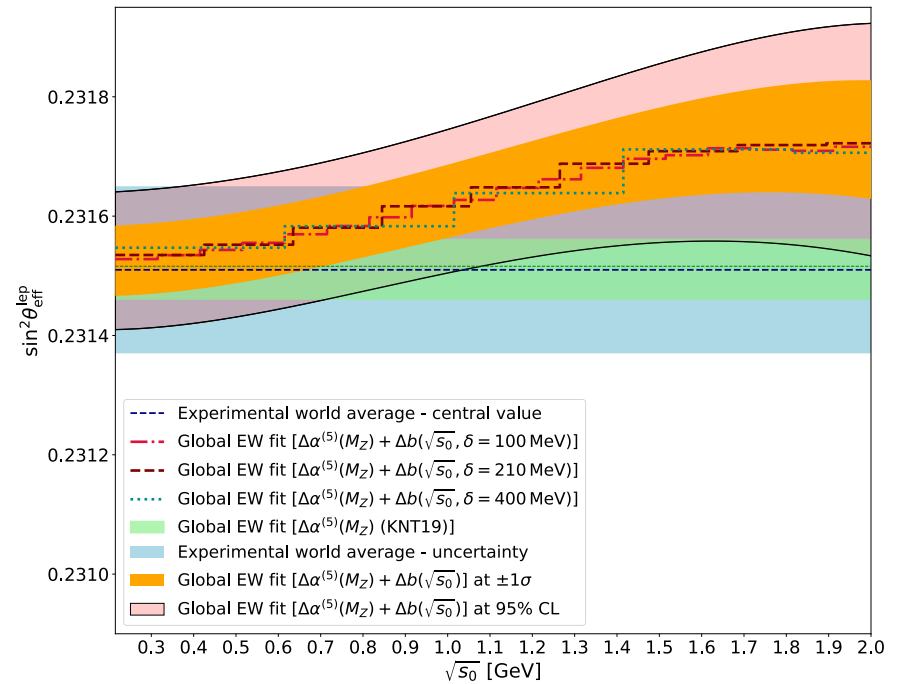
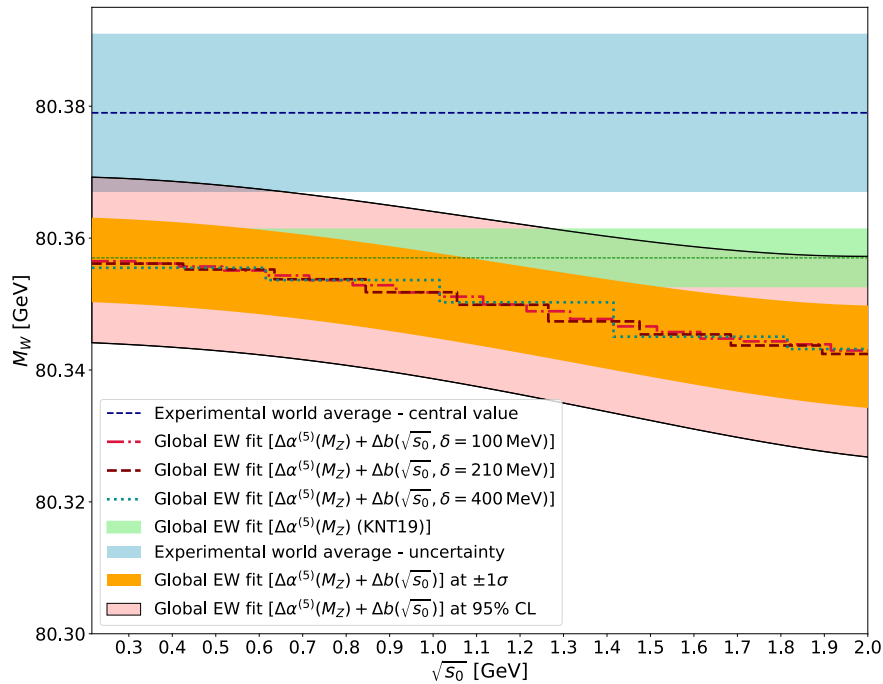
**Next MUonE theory workshop: MITP Mainz 01.03 - 05.03 2021 (hopefully)**



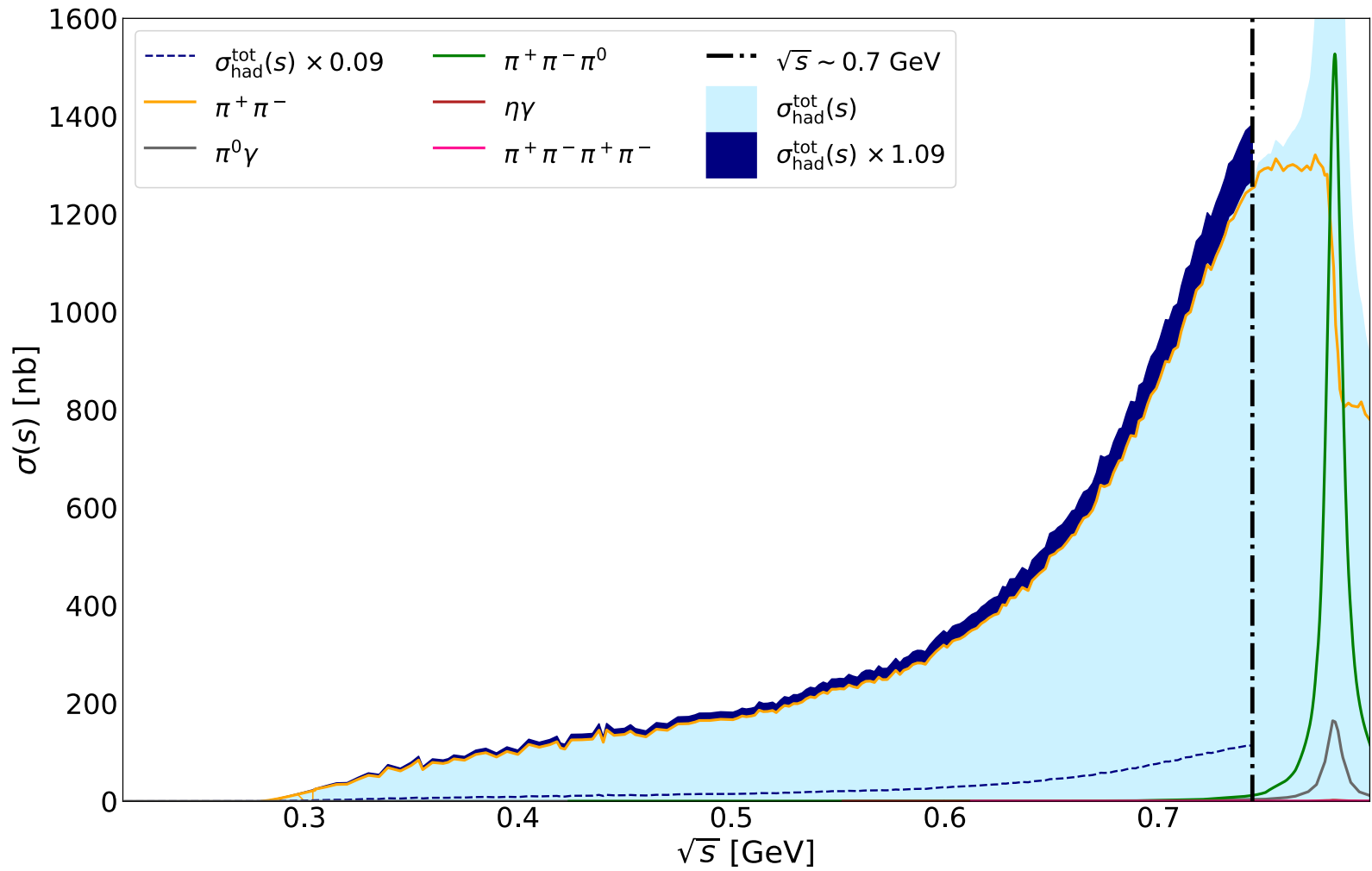
# Conclusions

- FNAL's first  $a_\mu$  result expected very soon with  $\sim$ BNL precision.
- Is present  $\Delta a_\mu$  due to missed contributions in the hadronic  $\sigma(s)$ ?
  - Shifts  $\Delta\sigma(s)$  to fix  $\Delta a_\mu$  conflict with the global EW fit above  $\sim 1$  GeV
  - Shifts below  $\sim 1$  GeV conflict with the quoted exp. error of  $\sigma(s)$ .
- Shifts  $\Delta\sigma(s)$  to fix  $\Delta a_\mu$  slightly increase the  $a_e$  tension ( $R_{e/\mu}$  ok).
- MUonE at CERN will provide a new independent spacelike determination of  $a_\mu^{\text{HLO}}$  alternative to the DR and lattice ones.

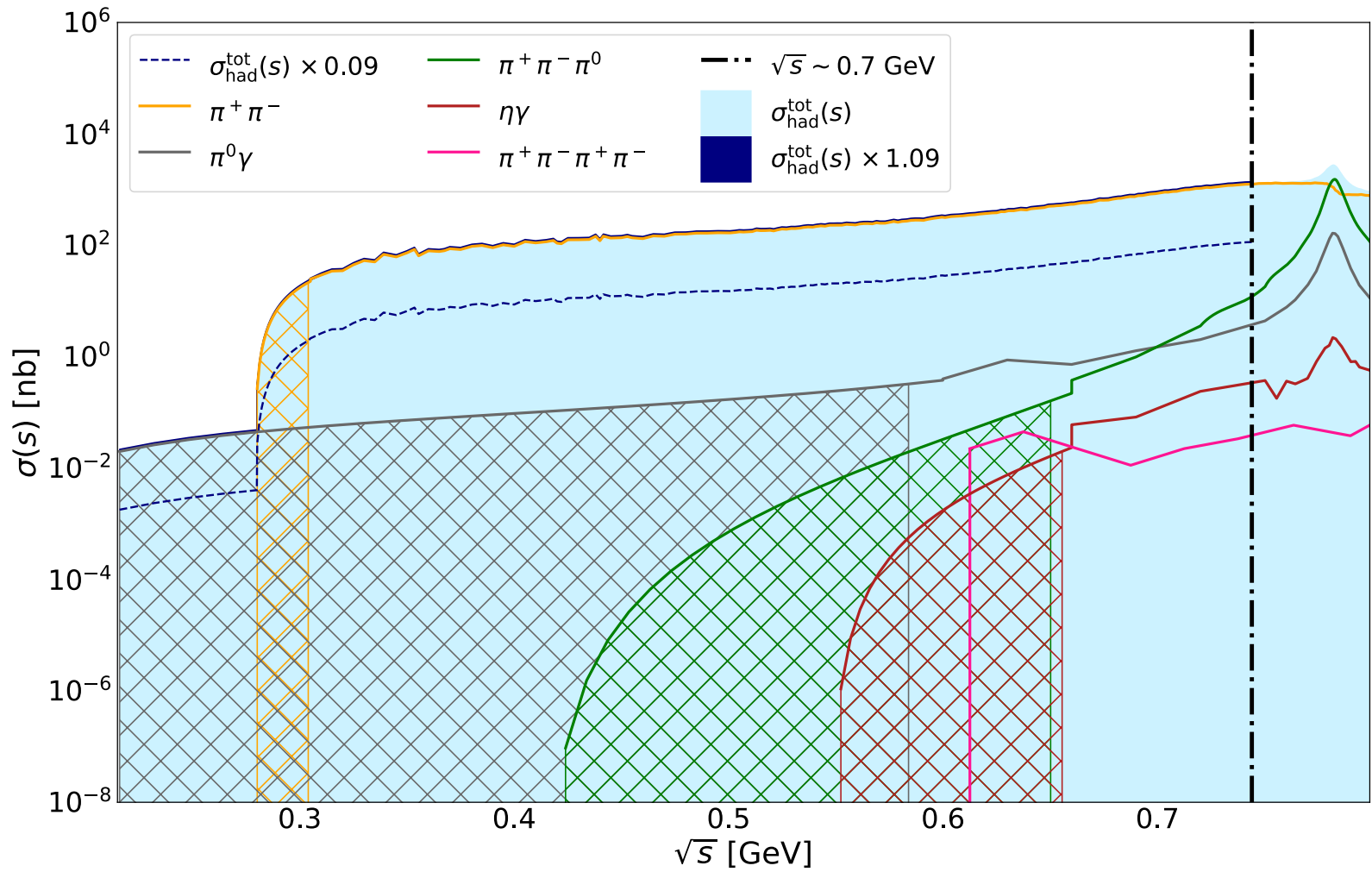
# Backup



Keshavarzi, Marciano, MP, Sirlin, PRD 2020



Keshavarzi, Marciano, MP, Sirlin, PRD 2020



Keshavarzi, Marciano, MP, Sirlin, PRD 2020